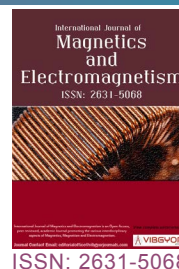


Angular Momentum Emission by a Rotating Dipole



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Abstract

A new calculation confirms the presence of spin radiation along the axis of rotation of a dipole. This is further proof of the need to introduce the spin tensor into classical electrodynamics, along with the energy-momentum tensor.

Keywords

Classical spin, Electrodynamics, Spin radiation

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Introduction

Circularly polarized electromagnetic radiation contains angular momentum in the form of the angular momentum density [1,2].

JH Poynting [2]: "If we put E for the energy in unit volume and G for the torque per unit area, we have $G = E\lambda / 2\pi$ ".

This means that such radiation is Weysenhoff's spin-fluid [3].

J Weysenhoff: "By spin-fluid we mean a fluid each element of which possesses besides energy and linear momentum also a certain amount of angular momentum, proportional - just as energy and the linear momentum - to the volume of the element".

This is recorded in textbooks [4,5]. Since Emma

Noether, this angular momentum has been described by the spin tensor density [6-8].

$$Y^{\lambda\mu\nu} = -2A^{[\lambda} \delta_{\alpha}^{\mu]} \frac{\partial L}{\partial(\partial_\nu A_\alpha)} = -2A^{[\lambda} F^{\mu]\nu} \quad (1)$$

Where $L = -F_{\mu\nu} F^{\mu\nu} / 4$ is the free electromagnetic field Lagrangian, A^λ is the vector potential, and $F_{\mu\nu}$ is the field-strength tensor. The local sense of a spin tensor is as follows. Y^{xyt} [J_s/m^3] is spin volume density, Y^{xyt} [J/m^2] is spin flux density, i.e. torque per unit area (cf. J. H. Poynting). The spin tensor is used in the publications [9-20]. However, the spin tensor is ignored in works expressing the common point of view, e.g. [21-25].

Besides spin, any electromagnetic field contains mass-energy and momentum, which are described by the energy-momentum tensor [26,27].

$$T^{\mu\nu} = -g^{\mu\lambda} F_{\lambda\alpha} F^{\nu\alpha} + g^{\mu\nu} F_{\alpha\beta} F^{\alpha\beta} / 4 \quad (2)$$

The local sense of the energy-momentum ten-

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tor is as follows. T^{xt} [N_ss/m³] is momentum volume density, T^{tx} [kg/m²s] is mass-energy flux density. It means, e.g., $dp^x = T^{xt}dV$ is the momentum in the volume dV .

Moment of momentum, e.g., $dL^{xy} = (xT^{yt} - yT^{xt})dV$ is the orbital angular momentum of the momentum contained in the volume dV . So, the total angular momentum possessed by the volume dV is

$$dJ^{ik} = dS^{ik} + dL^{ik} = (Y^{ikt} + 2r^{[i}T^{k]t})dV \quad (3)$$

The total torque per the area da_l , i.e. angular momentum flux, is

$$d\tau^{ik} = d\tau_s^{ik} + dL^{ik} / dt = (Y^{ikl} + 2r^{[i}T^{k]l})da_l \quad (4)$$

It is important that spin is not associated with a moment of a linear momentum, or even with a motion of matter. **Hehl** writes about spin of an electron [28]:

“The current density in Dirac’s theory can be split into a convective part and a polarization part. The polarization part is determined by the spin distribution of the electron field. It should lead to *no* energy flux in the rest system of the electron because the genuine spin ‘motion’ take place only within a region of the order of the Compton wavelength of the electron”.

Electromagnetic Field of a Rotating Dipole

Electromagnetic field of a rotating dipole \mathbf{p} is well known [27,29,30].

$$\mathbf{E} = \left[\frac{\omega^2(\mathbf{p}r^2 - (\mathbf{p}\mathbf{r})\mathbf{r})}{4\pi\epsilon_0c^2r^3} + \frac{i\omega(\mathbf{p}r^2 - 3(\mathbf{p}\mathbf{r})\mathbf{r})}{4\pi\epsilon_0cr^4} - \frac{(\mathbf{p}r^2 - 3(\mathbf{p}\mathbf{r})\mathbf{r})}{4\pi\epsilon_0r^5} \right] \exp(ikr - i\omega t) \quad (5)$$

$$\mathbf{H} = \left[\frac{\omega^2\mathbf{r} \times \mathbf{p}}{4\pi cr^2} + \frac{i\omega\mathbf{r} \times \mathbf{p}}{4\pi r^3} \right] \exp(ikr - i\omega t) \quad (6)$$

The first terms of (5), (6) are proportional to $1/r$ and so represent radiation. This radiation is of circular polarization in the direction of the rotational axis, z-axis (see **Figure 1** from [31]). Therefore this field contains the spin flux Y^{xyl} . We calculate this spin flux per sphere $r = Const$ in Section 3.

At the same time this radiation contains no orbital angular momentum flux per elements da_l of the sphere $r = Const$. $dL^{ik} / dt = 2r^{[i}T^{k]l}da_l = 0$. Really, the first terms fields \mathbf{E} & \mathbf{H} are orthogonal to each other and to the vector \mathbf{r} . So, in any point, we can enter lo-

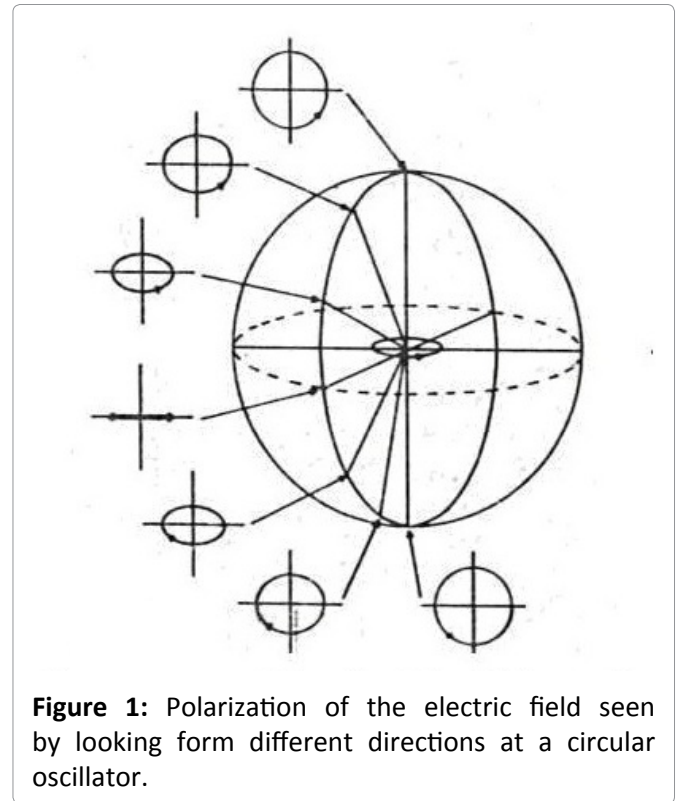


Figure 1: Polarization of the electric field seen by looking from different directions at a circular oscillator.

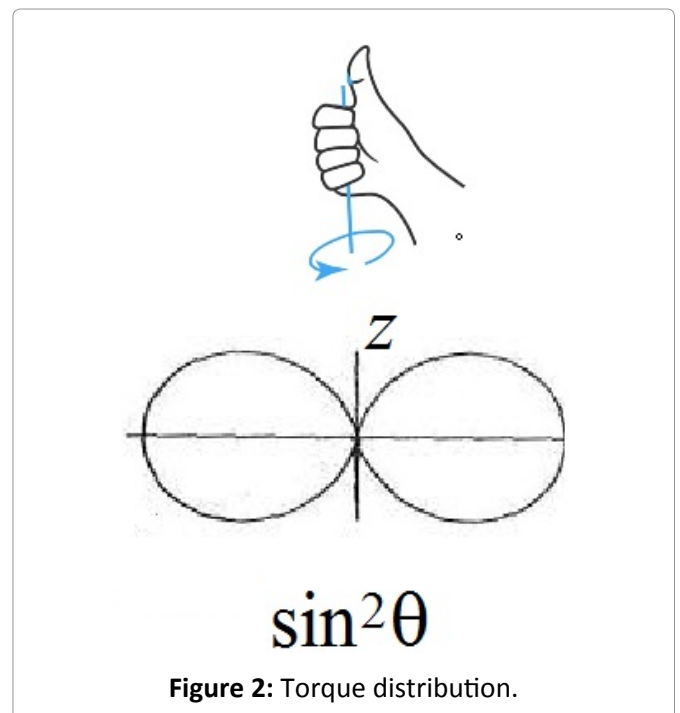


Figure 2: Torque distribution.

cal Cartesian coordinates such that $da_l = \{0, 0, da_z\}$, $\mathbf{E} = \{E_x, 0, 0\}$, $\mathbf{H} = \{0, H_y, 0\}$, $\mathbf{r} = \{0, 0, z\}$, i.e. $F_{tx}, F^{tx}, F_{xz}, F^{xz}$ are not equal to zero only. Using this coordinates we find according to (2): $T^{xz} = -g^{xx}F_{x\alpha}F^{z\alpha} = 0$, $T^{yz} = -g^{yy}F_{y\alpha}F^{z\alpha} = 0$. So the orbital angular momentum is not radiated.

The second terms field of (5), (6) contains the

orbital angular momentum flux, or torque, per the sphere $r = Const$. In Refs [32-37], spherical coordinates were used, and the angular distribution of the torque was obtained (see Figure 2):

$$dL^{ik} / dtd\Omega = \omega^3 p^2 \sin^2 \theta / 16\pi^2 \epsilon_0 c^3 \quad (7)$$

Where, $d\Omega = \sin \theta d\theta d\phi$. This torque is located in the neighborhood of the plane of rotation where the polarization is near linear. This torque is not radiated. This torque is like a static torque that someone can apply (Figure 2).

Spin Radiation by a Rotating Dipole

Spin radiated by the first terms field was calculated in [15] using the spin volume density Y^{xyt} on the assumption that this density is moving at the speed of light. Here the spin flux density Y^{xyt} is used. This is more naturally.

Using

$$\mathbf{E} = \frac{\omega^2(\mathbf{p}r^2 - (\mathbf{p}\mathbf{r})\mathbf{r})}{4\pi\epsilon_0 c^2 r^3} \exp(ikr - i\omega t), \quad \mathbf{H} = \frac{\omega^2 \mathbf{r} \times \mathbf{p}}{4\pi cr^2} \exp(ikr - i\omega t),$$

$$p_x = p, \quad p_y = ip \quad (8)$$

yields

$$E_x = F_{ix} = \frac{\omega^2 p(r^2 - x^2 - ixy)}{4\pi\epsilon_0 c^2 r^3},$$

$$E_y = F_{iy} = \frac{\omega^2 p(ir^2 - xy - iy^2)}{4\pi\epsilon_0 c^2 r^3},$$

$$E_z = F_{iz} = \frac{-\omega^2 p(zx + izy)}{4\pi\epsilon_0 c^2 r^3}, \quad (9)$$

$$H_x = F^{zy} = \frac{-i\omega^2 pz}{4\pi cr^2},$$

$$H_y = F^{xz} = \frac{\omega^2 pz}{4\pi cr^2},$$

$$H_z = F^{yx} = \frac{\omega^2 p(ix - y)}{4\pi cr^2} \quad (10)$$

Using $\mathbf{A} = -\int \mathbf{E} dt = -i\mathbf{E} / \omega$ yields

$$A_x = \frac{\omega p(-ir^2 + ix^2 - xy)}{4\pi\epsilon_0 c^2 r^3}, \quad A_y = \frac{\omega p(r^2 + ixy - y^2)}{4\pi\epsilon_0 c^2 r^3},$$

$$A_z = \frac{\omega p(izx - zy)}{4\pi\epsilon_0 c^2 r^3} \quad (11)$$

Accordingly to $Y^{\lambda\mu\nu} = -2A^{[\lambda} F^{\mu]\nu}$, we have

$$Y^{yx} = -\frac{\Re}{2} \{\bar{A}^x F^{yx}\} = \frac{\omega^3 z^2 x}{32\pi^2 \epsilon_0 c^3 r^5},$$

$$Y^{yy} = \frac{\Re}{2} \{\bar{A}^y F^{yy}\} = \frac{\omega^3 z^2 y}{32\pi^2 \epsilon_0 c^3 r^5},$$

$$Y^{yz} = -\frac{\Re}{2} \{\bar{A}^x F^{yz} - \bar{A}^y F^{xz}\} = \frac{\omega^3 (r^2 + z^2) z}{32\pi^2 \epsilon_0 c^3 r^5} \quad (12)$$

Because of $d\tau_S^{ik} = Y^{ikl} da_l$, we need the Cartesian coordinates of elements of the sphere $r = Const$, which spherical coordinates are $da_v = \{da_r = d\theta d\phi, da_\theta = 0, da_\phi = 0\}$. The transformation coefficients are; $\frac{\partial r}{\partial x} = \frac{x}{r}, \frac{\partial r}{\partial y} = \frac{y}{r}, \frac{\partial r}{\partial z} = \frac{z}{r}$, and $\sqrt{g} = r^2 \sin \theta$. So we have

$$da_l = \{da_x = x \sin \theta d\theta d\phi, da_y = y r \sin \theta d\theta d\phi, da_z = z r \sin \theta d\theta d\phi\},$$

and

$$d\tau_S^{xy} = Y^{xyl} da_l = Y^{yx} da_x + Y^{yy} da_y + Y^{yz} da_z$$

$$= \frac{\omega^3 p^2 (z^2 x^2 + z^2 y^2 + r^2 z^2 + z^4)}{32\pi^2 \epsilon_0 c^3 r^4} \sin \theta d\theta d\phi = \frac{\omega^3 p^2}{16\pi^2 \epsilon_0 c^3} \cos^2 \theta \sin \theta d\theta d\phi \quad (13)$$

This result, $d\tau_S^{xy} / d\Omega = \frac{\omega^3 p^2}{16\pi^2 \epsilon_0 c^3} \cos^2 \theta$, is coincided with Ref. [15]. The angular distribution of the spin radiation is represent in Figure 3.

Conclusion

A rotating electric dipole emits angular momen-

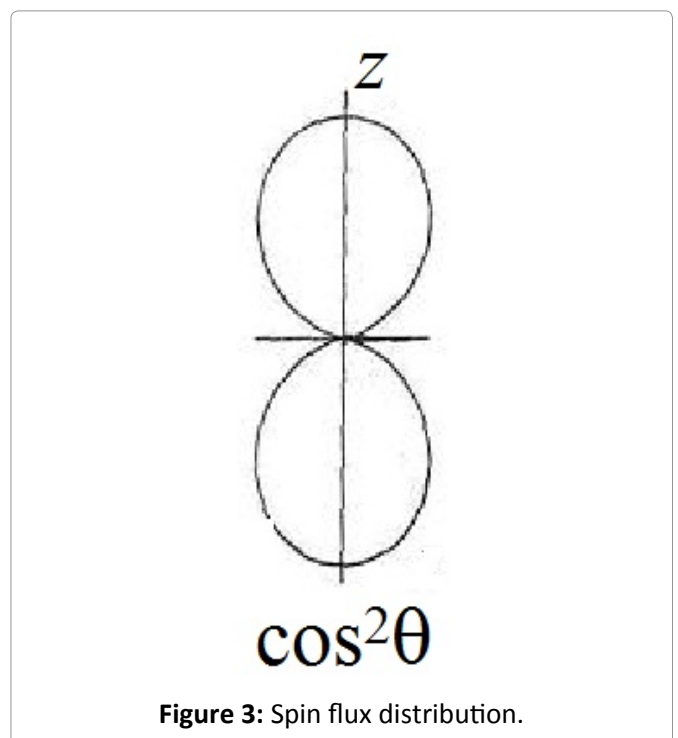


Figure 3: Spin flux distribution.

tum flux of two types: (i) Spin flux, which is directed mainly along the axis of rotation and determined by the spin tensor, and (ii) Orbital angular momentum flux determined by the energy-momentum tensor. The spin flux is not recognized by nowadays electrodynamics.

I am eternally grateful to Professor Robert Romer for the courageous publication of my question: "Does a plane wave really not carry spin?" [38] (was submitted on 07 October, 1999).

References

- Sadowsky A (1899) Acta et Comm. Imp. Universitatis Jurievensis, Russia, 7: 1-3.
- Poynting JH (1909) The wave motion of a revolving shaft, and a suggestion as to the angular momentum in a beam of circularly polarised light. Proc R Soc Lond A 82: 560-567.
- Weysenhoff J, Raabe A (1947) Relativistic dynamics of spin-fluids and spin-particles. Acta Phys Polon 9: 7-19.
- Crawford FS (1968) Waves: Berkley physics course. Berkeley, California, USA, 3.
- Feynman RP, Leighton RB, Sands M (1965) The Feynman lectures on physics. Addison-Wesley, London, 3: 17.
- Corson EM (1953) Introduction to tensors, spinors, and reativistic wave-equation. Hafner, USA, 71.
- Soper DE (2008) Classical field theory. Dover, USA, 114.
- Barut AO (1964) Electrodynamics and classical theory of particles and fields. Macmillan, USA, 102.
- Khrapko RI (2018) Absorption of spin by a conducting medium. AASCIT Journal of Physics 4: 59-63.
- Khrapko RI (2018) Absorption of angular momentum of a plane wave. Optik 154: 806-810.
- Khrapko R (2019) Unknown spin radiation. J Phys: Conf Ser 1172: 012055.
- Khrapko RI (2006) Origin of spin: Paradox of the classical Beth experiment. In: Smarandache F, Christian-to V, Yuhua Fu, Khrapko R, Hutchison J, Unfolding the Labyrinth: Open Problems in Mathematics, Physics, Astrophysics, and other areas of science. Hexis - Phoenix, 57-71.
- Khrapko RI (2008) Mechanical stresses produced by a light beam. J Modern Optics 55: 1487-1500.
- Khrapko RI (2017) Reflection of light from a moving mirror. Optik 136: 503-506.
- Khrapko RI (2019) Spin radiation from a rotating dipole. Optik 181: 1080-1084.
- Khrapko RI (2020) Radiation damping of a rotating dipole. Optik 203: 164021.
- Khrapko RI (2020) Absorption of spin of a plane circularly polarized wave. Optik 210: 164527.
- Khrapko RI (2002) True energy-momentum tensors are unique. Electrodynamics spin tensor is not zero. Moscow Aviation Institute, Russia.
- Khrapko RI (2001) Violation of the gauge equivalence. Moscow Aviation Institute, Russia.
- Khrapko RI (2005) Spin transmitted to the mirror when light is reflected. Moscow Aviation Institute, Russia.
- Andrews DL, Babiker M (2013) The angular momentum of light. Cambridge University Press, UK.
- Heitler W (1954) The quantum theory of radiation. Oxford University Press, Clarendon, UK, 401.
- Allen L, Padgett MJ (2002) Response to question #79. Does a plane wave carry spin angular momentum? Am J Phys 70: 567.
- Simmonds JW, Guttman MJ (1970) States, waves and photons. Addison-Wesley, Reading, USA.
- Ohanian HC (1986) What is spin? Amer J Phys 54: 500-505.
- Landau LD, Lifshitz EM (1975) The classical theory of fields. Pergamon, NY, USA.
- Jackson JD (1999) Classical electrodynamics. John Wiley, USA, 350.
- Hehl FW (1976) On the energy tensor of spinning massive matter in classical field theory and general relativity. Reports on Mathematical Physics 9: 55-82.
- Becker R (1982) Electromagnetic fields and interactions. Dover, NY, USA, 1: 284.
- Corney A (1977) Atomic and laser spectroscopy. Oxford University Press, Clarendon, UK, 36.
- Meyers RA (1987) Encyclopedie of physics science and technology. Academic Press, NY, USA, 2: 266.
- Sommerfeld A (1951) Atombau und Spektrallinien. Vieweg F, Braunschweio S, UK.
- Vul'fson KS (1987) Angular momentum of electromagnetic waves. Sov Phys Usp 30: 724-728.
- Barabanov AL(1993) Angular momentum in classical electrodynamics. Phys Usp 36: 1068-1074.
- Khrapko RI (2001) Spin of dipole radiation. Moscow Aviation Institute, Russia, 34635.
- Khrapko RI (2003) Radiation of spin by a rotator. Mathematical Physics Preprint Archive, 03-315.
- Khrapko RI (2012) Spin is not a moment of momentum. Moscow Aviation Institute, Russia, 28834.
- Khrapko RI (2001) Does plane wave not carry a spin? Amer J Phys 69: 405.